On some constructive aspects of an inverse problem in paleomagnetism

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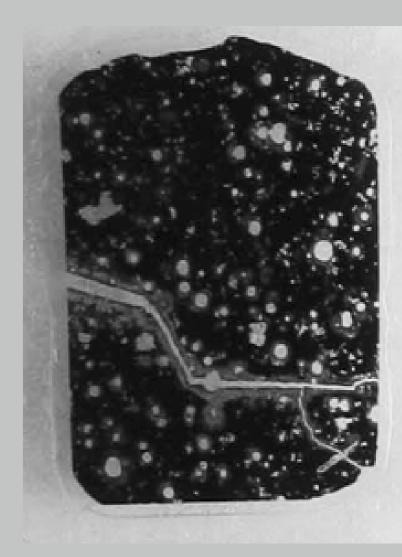
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Introduction



SQUID microscope (EAPS, MIT)

The process of extraction of relict magnetic information from georocks and meteorites is a challenging task. Due to the weak intensity of the field produced by the remanent magnetisation of a rock, the measurements have to be performed in direct vicinity of a sample and using highly sensitive magnetometric devices such as SQUID and QDM. Full reconstruction of the magnetic distribution (3D vector field) from partial measurements of only one component of the magnetic field is difficult and often replaced by estimation of the net magnetisation of a sample. We show some explicit asymptotic results for the latter problem, but observe that the lack of measured data and the presence of noise still remain an issue. We are thus motivated to consider the problem of stable extrapolation of magnetic field measurements.

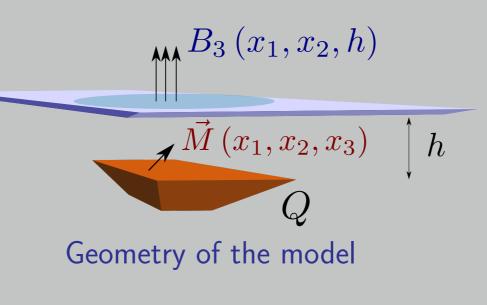


Slice of a magnetised sample (basalt)



1. Setting

• Unknown sample magnetisation: $\vec{\mathcal{M}}(\vec{x}) \equiv (\mathcal{M}_1, \mathcal{M}_2, \mathcal{M}_3) (\boldsymbol{x}, x_3),$ with supp $\vec{\mathcal{M}} \subset Q, \quad |Q| < \infty,$ $\boldsymbol{x} \equiv [x_1, x_2]^T, \quad \vec{x} \equiv [\boldsymbol{x}, x_3]^T.$



Measured component of the magnetic field:

$$B_{3}(\boldsymbol{x},h) = \frac{1}{4\pi} \frac{\partial}{\partial h} \iiint_{Q} \frac{\nabla \cdot \vec{\mathcal{M}}(\boldsymbol{t},\boldsymbol{t}_{3})}{\left(|\boldsymbol{x}-\boldsymbol{t}|^{2} + (h-t_{3})^{2}\right)^{1/2}} d^{3}\vec{t}, \quad \boldsymbol{x} \in \mathbb{R}^{2}$$
$$\begin{bmatrix} \vec{B} = -\nabla \Phi \quad \text{outside} \quad Q, \qquad \Delta \Phi = \nabla \cdot \vec{\mathcal{M}} \quad \text{in} \quad \mathbb{R}^{3} \end{bmatrix}$$

- ▶ Recovery *M*(*x*) ← *B*₃(*x*, *h*) is a severely <u>ill</u>-posed problem:
 ▷ The range of the mapping *∇* · *M* → *B*₃ is not closed.
 ▷ The mapping *M* → *∇* · *M* is not injective.
- Net magnetisation $\vec{m} := \iiint \vec{k} (\vec{x}) d^3 \vec{x}$ is unique and important.

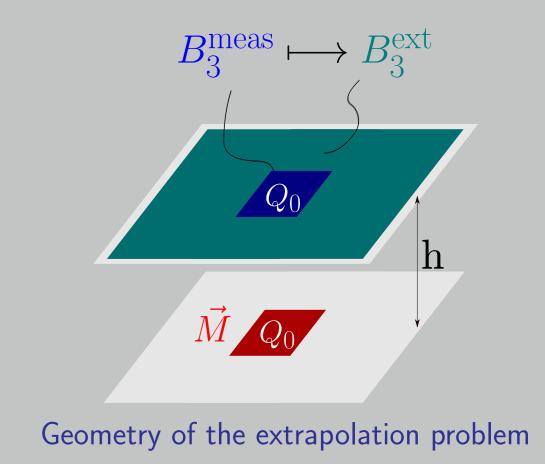
3. Overcoming the data limitation

Experimental limitations

- Measurements are available on a fairly small planar area.
- Only one component of the field \vec{B} is measured: B_3 .
- Away from sample, the data is heavily contaminated by <u>noise</u>.

Assumptions

- ► Planar regular magnetisation: $\mathcal{M}(\vec{x}) = \vec{M}(\boldsymbol{x}) \otimes \delta(x_3),$ supp $\vec{M} \subset Q_0 \subset \mathbb{R}^2,$ $\vec{M} \in [W^{1,2}(\mathbb{R}^2)]^2 \times L^2(\mathbb{R}^2).$
- \blacktriangleright B_3 is measured over Q_0 .



Extrapolation of measurements

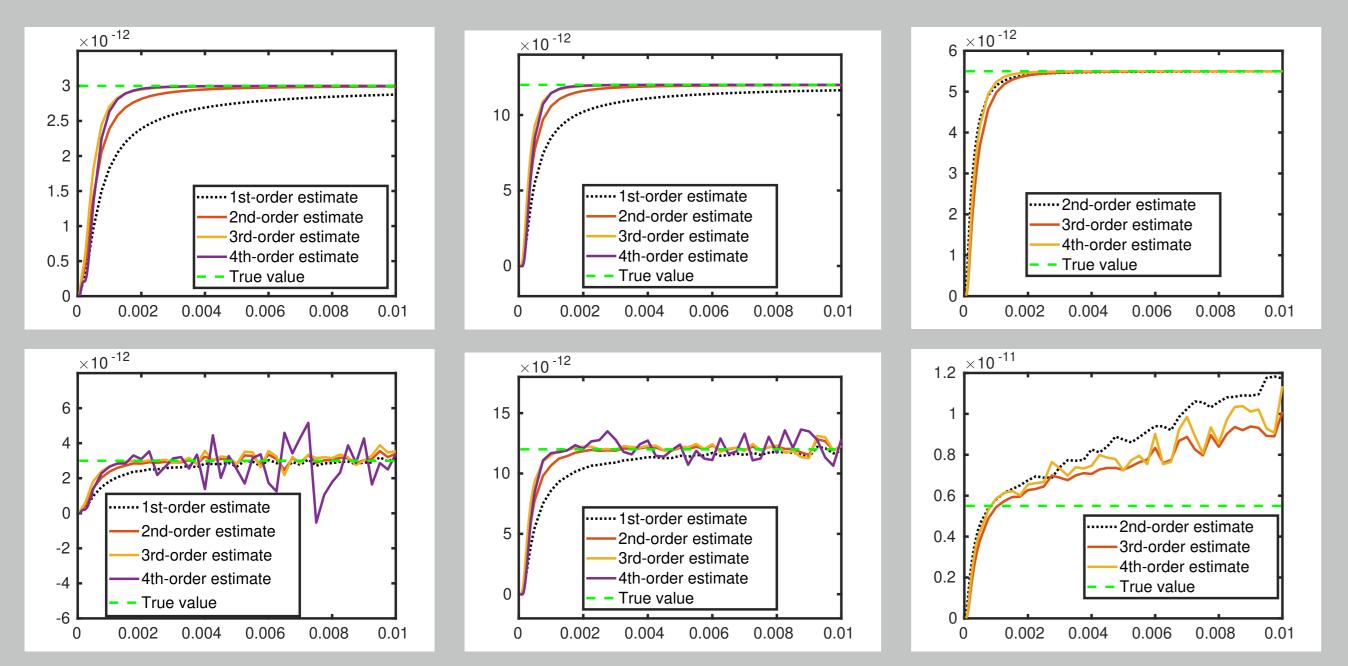
 \blacktriangleright We mainly aim to extrapolate B_3 on the plane:

2. Net magnetisation: asymptotic results

We can asymptotically estimate (explicitly and to an arbitrary order) $\vec{m} \leftarrow B_3(\cdot, h)|_{D_A}, \qquad D_A := \{ \boldsymbol{x} \in \mathbb{R}^2 : |\boldsymbol{x}| < A \}, \qquad A \gg 1.$

For example, 3rd-order estimates are given by

$$m_{j} = \frac{2}{5} \iint_{D_{A}} \left[5 + 24 \left(\frac{x_{j}}{A} \right)^{4} \right] x_{j} B_{3} (\boldsymbol{x}, h) \mathrm{d}^{2} \boldsymbol{x} + \mathcal{O} \left(\frac{1}{A^{3}} \right), \qquad j \in \{1, 2\}$$
$$m_{3} = \frac{A}{4} \iint_{D_{A}} \left[5 + 40 \left(\frac{x_{j}}{A} \right)^{4} - 128 \left(\frac{x_{j}}{A} \right)^{6} \right] B_{3} (\boldsymbol{x}, h) \mathrm{d}^{2} \boldsymbol{x} + \mathcal{O} \left(\frac{1}{A^{3}} \right).$$



 $B_3(\cdot,h)|_{Q_0} \coloneqq B_3^{\text{meas}} \longmapsto B_3^{\text{ext}} \coloneqq B_3(\cdot,h) \text{ on } \mathbb{R}^2.$

- ▶ B_1 , B_2 can be reconstructed from B_3^{ext} by suitable <u>bounded</u> integral operators. Moreover, we can obtain \vec{B} on the half–space $\{x_3 > h\}$.
- Extrapolation should at least be stable w.r.t. high-frequency noise.

4. Extrapolation method and results

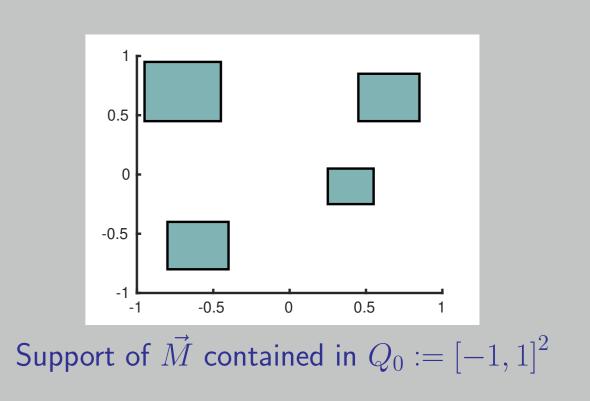
- 1. Fix $J, N \in \mathbb{N}_+$ sufficiently large.
- 2. Solve $K_{12} \star_{Q_0} \phi_j = \mu_j \phi_j$ on Q_0 . $\Rightarrow J \text{ largest } (\mu_j)_{j=1}^J \subset \mathbb{R}_+$ and $(\phi_j)_{j=1}^J \subset L^2_{\mathbb{R}}(Q_0), \|\phi_j\| = 1.$
- 3. Compute on $Q_0 \times Q_0$ $S_J(\mathbf{x}, \mathbf{t}) := \sum_{j=1}^J \frac{\langle K_{12}(\cdot - \mathbf{x}), \phi_j \rangle}{\mu_j} \phi_j(\mathbf{t}).$
- 4. Solve on Q_0 $\begin{cases}
 K_{12} \star_{Q_0} \varphi_{12}^n + K_3 \star_{Q_0} \varphi_3^n = \lambda_n \varphi_{12}^n, \\
 K_{12} \star_{Q_0} \varphi_{12}^n + \mathcal{S}_J K_3 \star_{Q_0} \varphi_3^n = \lambda_n \varphi_3^n.
 \end{cases}$ $\Rightarrow N \text{ largest in modulus } (\lambda_n)_{n=1}^N \subset \mathbb{R}$ and $\left([\varphi_{12}^n, \varphi_3^n]^T \right)_{n=1}^N \subset [L^2_{\mathbb{R}}(Q_0)]^2, \\
 \|\varphi_{12}^n\|^2 + \|\varphi_3^n\|^2 = 1.
 \end{cases}$

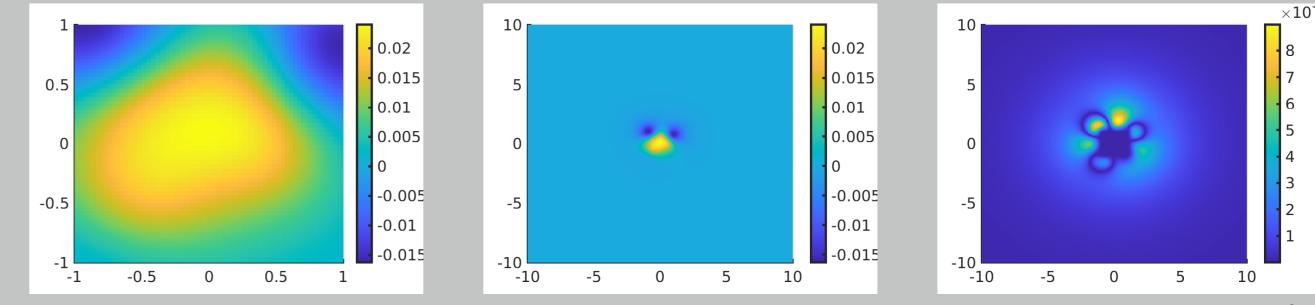
 $B_{3} = K_{12} \star_{Q_{0}} D_{M_{12}} + K_{3} \star_{Q_{0}} M_{3}$ Notations $K_{12}(\mathbf{x}) := -\frac{h}{4\pi (|\mathbf{x}|^{2} + h^{2})^{3/2}},$ $K_{3}(\mathbf{x}) := \partial_{h} K_{12}(\mathbf{x}),$ $D_{M_{12}}(\mathbf{x}) := \partial_{1} M_{1}(\mathbf{x}) + \partial_{2} M_{2}(\mathbf{x}),$ $(K \star_{Q_{0}} f)(\mathbf{x}) \equiv \iint_{Q_{0}} K(\mathbf{x} - \mathbf{t}) f(\mathbf{t}) d^{2}\mathbf{t},$ $(\mathcal{S}_{J}f)(\mathbf{x}) := \iint_{Q_{0}} S_{J}(\mathbf{x}, \mathbf{t}) f(\mathbf{t}) d^{2}\mathbf{t},$ $\langle f, g \rangle \equiv \iint_{Q_{0}} f(\mathbf{t}) g(\mathbf{t}) d^{2}\mathbf{t}.$

Recovery of m_1 , m_2 , m_3 (left to right) for the case of pure (top) and contaminated (bottom) field

Some references

- L. Baratchart, D.P. Hardin, E.A. Lima, E.B. Saff, B.P. Weiss, "Characterizing kernels of operators related to thin-plate magnetizations via generalizations of Hodge decompositions", *Inverse Problems 29*, 2013.
- D. Ponomarev, "Magnetisation moment of a bounded 3D sample: asymptotic recovery from planar measurements on a large disk using Fourier analysis", arXiv:2205.14776.
- 5. Construct $B_3^{\text{ext}}(\mathbf{x}) = \sum_{n=1}^{N} b_n \widetilde{\varphi}_{12}^n(\mathbf{x}),$ $b_n := \langle B_3^{\text{meas}}, \varphi_{12}^n \rangle + \langle S_J B_3^{\text{meas}}, \varphi_3^n \rangle,$ $\widetilde{\varphi}_{12}^n(\mathbf{x}) := \frac{1}{\lambda_n} \iint_{Q_0} [K_{12}(\mathbf{x} - \mathbf{t}) \varphi_{12}^n(\mathbf{t}) + K_3(\mathbf{x} - \mathbf{t}) \varphi_3^n(\mathbf{t})] d^2 \mathbf{t}, \quad \mathbf{x} \in \mathbb{R}^2.$





Original $B_3^{
m meas}$ (left), extrapolated field $B_3^{
m ext}$ (middle) and extrapolation error (right) on $[-10, 10]^2$

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